

# KANTBP 3.0: New Version of a Program for Computing Energy Levels, Reflection and Transmission Matrices, and Corresponding Wave Functions in the Coupled-Channel Adiabatic Approach

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Brief description of a FORTRAN 77 program for calculating energy values, reflection and transmission matrices, and corresponding wave functions in a coupled-channel approximation of the adiabatic approach is presented. In this approach, a multidimensional Schrödinger equation is reduced to a system of the coupled second-order ordinary differential equations on a finite interval with the homogeneous boundary conditions of the third type at the left- and right-boundary points for continuous spectrum problem, or a set of first, second and third type boundary conditions for discrete spectrum problem. The resulting system of these equations containing the potential matrix elements and first-derivative coupling terms is solved using high-order accuracy approximations of the finite element method.

**Key words and phrases:** boundary value problem, multichannel scattering problem, finite element method, Kantorovich method.

## 1. Introduction

In this work we present a brief description of a KANTBP3 program for calculating with a required accuracy approximate eigensolutions of the continuum spectrum for systems of coupled differential equations on finite intervals of the variable  $z \in [z_{\min}, z_{\max}]$  using a general homogeneous boundary condition of the third-type [1]. The third-type boundary conditions are formulated for problems under consideration by using known asymptotics for a set of linear independent asymptotic regular and irregular solutions in the open channels, and a set of linear independent regular asymptotic solutions in the closed channels, respectively [2]. These problems are solved by the finite element method [3, 4]. This approach can be used in calculations of effects of electron screening on low-energy fusion cross sections, channeling processes, threshold phenomena in the formation and ionization of (anti)hydrogen-like atoms and ions in magnetic traps, scattering problem for quantum dots and quantum wires in magnetic field, potential scattering with confinement potentials, penetration through a two-dimensional fission barrier, tunneling from false vacuum of two interacted particles and three-dimensional tunneling of a diatomic molecule incident upon a potential barrier [2, 5].

## 2. Statement of the Problem

In the Kantorovich method or close-coupling adiabatic approach, the multidimensional Schrödinger equation is reduced to a finite set of  $N$  ordinary second-order differential equations on the finite interval  $[z_{\min}, z_{\max}]$  for the partial solution

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$$\boldsymbol{\chi}^{(j)}(z) = \left( \chi_1^{(j)}(z), \dots, \chi_N^{(j)}(z) \right)^T,$$

$$(\mathbf{L} - 2E\mathbf{I})\boldsymbol{\chi}^{(j)}(z) = 0, \quad \mathbf{L} = -\mathbf{I} \frac{1}{z^{d-1}} \frac{d}{dz} z^{d-1} \frac{d}{dz} + \mathbf{V}(z) + \mathbf{Q}(z) \frac{d}{dz} + \frac{1}{z^{d-1}} \frac{d z^{d-1}}{dz} \mathbf{Q}(z). \quad (1)$$

Here  $\mathbf{I}$ ,  $\mathbf{V}(z)$  and  $\mathbf{Q}(z)$  are the unit, symmetric and antisymmetric  $N \times N$  matrices, respectively. We assume that  $\mathbf{V}(z)$  and  $\mathbf{Q}(z)$  matrices have the following asymptotic behaviour at large  $z = z_{\pm} \rightarrow \pm\infty$

$$V_{ij}(z_{\pm}) = \left( \epsilon_j + \frac{2Z_j^{\pm}}{z_{\pm}} \right) \delta_{ij} + \sum_{l=2} \frac{v_{ij}^{(l,\pm)}}{z_{\pm}^l}, \quad Q_{ij}(z_{\pm}) = \sum_{l=1} \frac{q_{ij}^{(l,\pm)}}{z_{\pm}^l}, \quad (2)$$

where  $\epsilon_1 \leq \dots \leq \epsilon_N$  are the threshold energy values.

In the present work, scattering problem is solved using the boundary conditions at  $d = 1$ ,  $z = z_{\min}$  and  $z = z_{\max}$ :

$$\left. \frac{d\boldsymbol{\Phi}(z)}{dz} \right|_{z=z_{\min}} = \mathcal{R}(z_{\min})\boldsymbol{\Phi}(z_{\min}), \quad \left. \frac{d\boldsymbol{\Phi}(z)}{dz} \right|_{z=z_{\max}} = \mathcal{R}(z_{\max})\boldsymbol{\Phi}(z_{\max}), \quad (3)$$

where  $\mathcal{R}(z)$  is a unknown  $N \times N$  matrix-function,  $\boldsymbol{\Phi}(z) = \{\boldsymbol{\chi}^{(j)}(z)\}_{j=1}^{N_o}$  is the required  $N \times N_o$  matrix-solution and  $N_o$  is the number of open channels,  $N_o = \max_{2E \geq \epsilon_j} j \leq N$ .

From this we obtain the quadratic functional at  $d = 1$  (similar to Eq. (5) in [3])

$$\begin{aligned} \Xi(\boldsymbol{\Phi}, E, z_{\min}, z_{\max}) &\equiv \int_{z_{\min}}^{z_{\max}} \boldsymbol{\Phi}^T(z) (\mathbf{L} - 2E\mathbf{I}) \boldsymbol{\Phi}(z) dz = \mathbf{\Pi}(\boldsymbol{\Phi}, E, z_{\min}, z_{\max}) - \\ &- \boldsymbol{\Phi}^T(z_{\max}) \mathbf{G}(z_{\max}) \boldsymbol{\Phi}(z_{\max}) + \boldsymbol{\Phi}^T(z_{\min}) \mathbf{G}(z_{\min}) \boldsymbol{\Phi}(z_{\min}), \end{aligned} \quad (4)$$

where  $\mathbf{\Pi}(\boldsymbol{\Phi}, E, z_{\min}, z_{\max})$  is the symmetric functional

$$\begin{aligned} \mathbf{\Pi}(\boldsymbol{\Phi}, E, z_{\min}, z_{\max}) &= \int_{z_{\min}}^{z_{\max}} \left[ \frac{d\boldsymbol{\Phi}^T(z)}{dz} \frac{d\boldsymbol{\Phi}(z)}{dz} + \boldsymbol{\Phi}^T(z) \mathbf{V}(z) \boldsymbol{\Phi}(z) + \right. \\ &\left. + \boldsymbol{\Phi}^T(z) \mathbf{Q}(z) \frac{d\boldsymbol{\Phi}(z)}{dz} - \frac{d\boldsymbol{\Phi}(z)^T}{dz} \mathbf{Q}(z) \boldsymbol{\Phi}(z) - 2E \boldsymbol{\Phi}^T(z) \boldsymbol{\Phi}(z) \right] dz, \end{aligned} \quad (5)$$

and  $\mathbf{G}(z) = \mathcal{R}(z) - \mathbf{Q}(z)$  is the  $N \times N$  matrix-function which should be symmetric according to the conventual  $\mathbf{R}$ -matrix theory.

### 3. The Physical Scattering Asymptotic Forms

Matrix-solution  $\boldsymbol{\Phi}_v(z) = \boldsymbol{\Phi}(z)$  describing the incidence of the particle and its scattering, which has the asymptotic form "incident wave + outgoing waves" is

$$\boldsymbol{\Phi}_v(z \rightarrow \pm\infty) = \begin{cases} \begin{cases} \mathbf{X}^{(+)}(z) \mathbf{T}_v, & z > 0, \\ \mathbf{X}^{(+)}(z) + \mathbf{X}^{(-)}(z) \mathbf{R}_v, & z < 0, \end{cases} & v = \rightarrow, \\ \begin{cases} \mathbf{X}^{(-)}(z) + \mathbf{X}^{(+)}(z) \mathbf{R}_v, & z > 0, \\ \mathbf{X}^{(-)}(z) \mathbf{T}_v, & z < 0, \end{cases} & v = \leftarrow, \end{cases} \quad (6)$$

where  $\mathbf{R}_v$  and  $\mathbf{T}_v$  are the reflection and transmission  $N_o \times N_o$  matrices,  $v = \rightarrow$  and  $v = \leftarrow$  denote the initial direction of the particle motion along the  $z$  axis. Here the leading term of the asymptotic rectangle-matrix functions  $\mathbf{X}^{(\pm)}(z)$  has the form [2]

$$X_{ij}^{(\pm)}(z) \rightarrow p_j^{-1/2} \exp\left(\pm i\left(p_j z - \frac{Z_j}{p_j} \ln(2p_j|z|)\right)\right) \delta_{ij}, \quad (7)$$

$$p_j = \sqrt{2E - \epsilon_j} \quad i = 1, \dots, N, \quad j = 1, \dots, N_o,$$

where  $Z_j = Z_j^+$  at  $z > 0$  and  $Z_j = Z_j^-$  at  $z < 0$ .

The matrix-solution  $\Phi_v(z, E)$  is normalized by

$$\int_{-\infty}^{\infty} \Phi_{v'}^\dagger(z, E') \Phi_v(z, E) dz = 2\pi \delta(E' - E) \delta_{v'v} \mathbf{I}_{oo}, \quad (8)$$

where  $\mathbf{I}_{oo}$  is the unit  $N_o \times N_o$  matrix. Let us rewrite Eq. (6) in the matrix form at  $z_+ \rightarrow +\infty$  and  $z_- \rightarrow -\infty$  as

$$\begin{pmatrix} \Phi_{\rightarrow}(z_+) & \Phi_{\leftarrow}(z_+) \\ \Phi_{\rightarrow}(z_-) & \Phi_{\leftarrow}(z_-) \end{pmatrix} = \begin{pmatrix} \mathbf{0} & \mathbf{X}^{(-)}(z_+) \\ \mathbf{X}^{(+)}(z_-) & \mathbf{0} \end{pmatrix} + \begin{pmatrix} \mathbf{0} & \mathbf{X}^{(+)}(z_+) \\ \mathbf{X}^{(-)}(z_-) & \mathbf{0} \end{pmatrix} \mathbf{S}, \quad (9)$$

where the scattering matrix  $\mathbf{S}$

$$\mathbf{S} = \begin{pmatrix} \mathbf{R}_{\rightarrow} & \mathbf{T}_{\leftarrow} \\ \mathbf{T}_{\rightarrow} & \mathbf{R}_{\leftarrow} \end{pmatrix} \quad (10)$$

is composed of the reflection and transmission matrices.

In addition, it should be noted that functions  $\mathbf{X}^{(\pm)}(z)$  satisfy relations

$$\mathbf{Wr}(\mathbf{Q}(z); \mathbf{X}^{(\mp)}(z), \mathbf{X}^{(\pm)}(z)) = \pm 2i \mathbf{I}_{oo}, \quad \mathbf{Wr}(\mathbf{Q}(z); \mathbf{X}^{(\pm)}(z), \mathbf{X}^{(\pm)}(z)) = \mathbf{0}, \quad (11)$$

where

$$\mathbf{Wr}(\bullet; \mathbf{a}(z), \mathbf{b}(z)) = \mathbf{a}^T(z) \left( \frac{d\mathbf{b}(z)}{dz} - \bullet \mathbf{b}(z) \right) - \left( \frac{d\mathbf{a}(z)}{dz} - \bullet \mathbf{a}(z) \right)^T \mathbf{b}(z). \quad (12)$$

This Wronskian is used to estimate a desirable accuracy of the above expansion.

Note, using a wronskian, we obtain the following properties of the reflection and transmission matrices:

$$\begin{aligned} \mathbf{T}_{\rightarrow}^\dagger \mathbf{T}_{\rightarrow} + \mathbf{R}_{\rightarrow}^\dagger \mathbf{R}_{\rightarrow} &= \mathbf{I}_{oo} = \mathbf{T}_{\leftarrow}^\dagger \mathbf{T}_{\leftarrow} + \mathbf{R}_{\leftarrow}^\dagger \mathbf{R}_{\leftarrow}, \\ \mathbf{T}_{\rightarrow}^\dagger \mathbf{R}_{\leftarrow} + \mathbf{R}_{\rightarrow}^\dagger \mathbf{T}_{\leftarrow} &= \mathbf{0} = \mathbf{R}_{\leftarrow}^\dagger \mathbf{T}_{\rightarrow} + \mathbf{T}_{\leftarrow}^\dagger \mathbf{R}_{\rightarrow}, \\ \mathbf{T}_{\rightarrow}^T &= \mathbf{T}_{\leftarrow}, \quad \mathbf{R}_{\rightarrow}^T = \mathbf{R}_{\rightarrow}, \quad \mathbf{R}_{\leftarrow}^T = \mathbf{R}_{\leftarrow}. \end{aligned} \quad (13)$$

This means that the scattering matrix (10) is symmetric and unitary.

#### 4. Test Desk

We consider the boundary problem (1)–(3) with parameters  $d = 1$ ,  $\hat{Z}_1 = \hat{Z}_2 = 0.1$ ,  $m_1 = 1$ ,  $m_2 = 3$ ,  $s = 8$ ,  $\bar{x}_{\min} = 0.1$ . This problem is followed from Kantorovich expansion of the 2D BVP described the tunneling problem of transmission of two ions

through repulsive barrier (for details, see [2])

$$\left( -\frac{\partial^2}{\partial y^2} - \frac{\partial^2}{\partial x^2} + x^2 + U_1(x_1) + U_2(x_2) - \mathcal{E} \right) \Psi(x, y) = 0, \quad (14)$$

where  $U_i(x_i) = 2\hat{Z}_i / \sqrt{|x_i|^s + \bar{x}_{\min}^s}$  are Coulomb-like barrier potentials,  $x_1 = s_2y + s_1x$  and  $x_2 = s_2y - s_3x$  are Jacobi coordinates with  $s_1 = m_1/M$ ,  $s_3 = m_2/M$ ,  $s_2 = \sqrt{m_1m_2}/M$ ,  $s_2 = \sqrt{\frac{\mu}{M}}$ ,  $M = m_1 + m_2$ .

The required asymptotics of regular and irregular solutions given in [2]. The following values of numerical parameters and characters described in [1] have been used in the test run via the supplied input file SQRTBT.INP

```
&PARAS TITLE=' REFLECTION AND TRANSMISSION MATRICES ',
      IPTYPE=1,NROOT=1,MDIM=4, IDIM=1,NPOL=4,
      SHIFT= 4D0, IPRINT=1, IPRSTP=120,
      NMESH=7, RMESH=-25D0, 100D0, -6D0, 100D0, 6D0, 100D0, 25D0,
      NDIR=1, NDIL=4, NMDIL=0, THRSHL= 1.D0, 3D0, 5D0, 7D0, IBOUND=8,
      FNOUT='KANTBP.LPR', IOUT=7, POTEN='ODPEVP.PTN', IOUP=10,
      FMATR='KANTBP.MAT', IOUM=11, EVWFN='KANTBP.WFN', IOUF=0
&END
```

Boundary problem (14) and the corresponding matrix elements  $\mathbf{V}(y)$ ,  $\mathbf{Q}(y)$  have been solved by the ODPEVP program [6] on grids  $\Omega_x\{x_{\min}, x_{\max}\} = \{-x_{\min}(64)x_{\max}\}$  with accuracy  $\epsilon_{ps} = 10^{-10}$ . Boundary points are  $x_{\max} = -x_{\min} = 8.1$ . All calculation details of this problem were written into file ODPEVP.LPR.

## TEST RUN OUTPUT

PROBLEM: REFLECTION AND TRANSMISSION MATRICES  
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### CONTROL INFORMATION

```
-----
NUMBER OF DIFFERENTIAL EQUATIONS. . . . . (MDIM ) =    4
NUMBER OF FINITE ELEMENTS . . . . . (NELEM ) =   300
NUMBER OF GRID POINTS . . . . . (NGRID ) =  1201
ORDER OF SHAPE FUNCTIONS. . . . . (NPOL ) =    4
ORDER OF GAUSS-LEGENDRE QUADRATURE. . . . (NGQ ) =    5
DIMENSION OF ENVELOPE SPACE . . . . . (IDIM ) =    1
BOUNDARY CONDITION CODE . . . . . (IBOUND) =    8
DOUBLE ENERGY SPECTRUM. . . . . (SHIFT ) =  4.00000
```

SUBDIVISION OF RHO-REGION ON THE FINITE-ELEMENT GROUPS:

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NO OF GROUP	NUMBER OF ELEMENTS	BEGIN OF INTERVAL	LENGTH OF ELEMENT	GRID STEP	END OF INTERVAL
1	100	-25.000	0.19000	0.04750	-6.000
2	100	-6.000	0.12000	0.03000	6.000
3	100	6.000	0.19000	0.04750	25.000

### TOTAL SYSTEM DATA

```
-----
TOTAL NUMBER OF ALGEBRAIC EQUATIONS. . . . (NN ) =   4804
TOTAL NUMBER OF MATRIX ELEMENTS. . . . . (NWK ) =  60010
MAXIMUM HALF BANDWIDTH . . . . . (MK ) =    20
MEAN HALF BANDWIDTH . . . . . (MMK ) =    12
```

NDIM, MDIM= 4 4

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CALCULATION OF WAVE FUNCTION WITH DIRECTION <--

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NUMBER OF OPEN CHANNELS. . . . . (NOPEN) = 2  
 VALUE OF I-TH MOMENTUM . . . . . (I, QR ) = 1 0.1732E+01  
 VALUE OF I-TH MOMENTUM . . . . . (I, QR ) = 2 0.1000E+01

I M P A R T: W R O N S K I A N

-----  
 -2.00000 -0.168196E-08  
 -0.168196E-08 -2.00000

\*\*\*\*\*

R E P A R T: R R M A T R I X

-----  
 -0.194759 -0.590855E-03  
 -0.590855E-03 -0.485377E-01

\*\*\*\*\*

I M P A R T: R R M A T R I X

-----  
 -0.124681 0.172716  
 0.172716 0.931470

\*\*\*\*\*

R E P A R T: T T M A T R I X

-----  
 0.600459 -0.317924E-01  
 0.317924E-01 -0.276468

\*\*\*\*\*

I M P A R T: T T M A T R I X

-----  
 -0.729781 0.150166  
 -0.150166 0.134581E-01

\*\*\*\*\*

Z R E P A R T: F U N C T I O N S

-----  
 -25.0000 0.6664D+00 -0.1165D+00 0.1531D+00 -0.1120D+00 0.7601D-06 0.8680D-05 0.2445D-07 0.4751D-06  
 -13.6000 0.6802D+00 -0.7978D-01 0.4223D-01 0.2431D+00 -0.2701D-04 0.3867D-04 -0.2948D-05 0.3128D-05  
 -3.6000 0.1490D-01 -0.5461D-01 -0.3718D-01 -0.2780D+00 -0.9230D-02 0.2404D-02 -0.1299D-02 -0.2425D-03  
 0.0000 -0.8416D+00 0.7861D-01 0.9335D-02 0.4446D+00 0.5115D-01 -0.1732D-01 -0.2247D-02 -0.6850D-02  
 3.6000 -0.4115D+00 -0.6691D-01 0.8351D-01 0.1351D+01 -0.4630D-02 0.2048D-01 -0.2890D-04 -0.9308D-04  
 13.6000 0.5769D+00 -0.6829D-01 -0.8088D-01 -0.1298D+01 -0.3777D-04 -0.5932D-04 0.3999D-05 0.9632D-05  
 25.0000 0.2716D+00 -0.1259D+00 -0.1631D+00 -0.5370D+00 0.1506D-05 0.4406D-04 -0.6915D-07 -0.2284D-05

\*\*\*\*\*

Z I M P A R T: F U N C T I O N S

-----  
 -25.0000 0.2735D+00 0.1055D-01 -0.2391D-01 -0.2560D+00 0.6563D-05 -0.4645D-05 0.3403D-06 -0.2162D-06  
 -13.6000 -0.2428D+00 0.8603D-01 0.1506D+00 -0.1425D+00 0.2784D-04 0.5248D-04 0.2187D-05 0.5597D-05  
 -3.6000 0.7372D+00 -0.1083D+00 -0.1518D+00 0.1221D+00 0.1107D-02 -0.5541D-02 -0.1592D-02 -0.3799D-03  
 0.0000 0.5262D+00 -0.1487D+00 -0.1846D-01 0.6235D+00 -0.3508D-01 -0.4223D-03 -0.5965D-02 -0.9388D-02  
 3.6000 -0.5284D+00 -0.8131D-01 0.1780D+00 0.1380D+01 0.1289D-01 0.1938D-01 -0.2320D-02 0.4662D-04  
 13.6000 -0.5507D+00 0.1129D+00 -0.1559D+00 -0.1335D+01 0.6059D-05 -0.8405D-04 -0.3894D-06 0.1222D-04  
 25.0000 -0.8982D+00 0.3837D-01 0.6149D-01 -0.6546D+00 0.5103D-05 0.4498D-04 -0.2851D-06 -0.2320D-05

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## CALCULATION OF WAVE FUNCTION WITH DIRECTION --&gt;

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NUMBER OF OPEN CHANNELS. . . . . (NOPEN) = 2  
 VALUE OF I-TH MOMENTUM . . . . . (I, QR ) = 1 0.1732E+01  
 VALUE OF I-TH MOMENTUM . . . . . (I, QR ) = 2 0.1000E+01

## I M P A R T: W R O N S K I A N

-----  
 2.00000      -.168196E-08  
 -.168196E-08    2.00000

\*\*\*\*\*

## R E P A R T: R R M A T R I X

-----  
 -.194759      0.590855E-03  
 0.590855E-03   -.485377E-01

\*\*\*\*\*

## I M P A R T: R R M A T R I X

-----  
 -.124681      -.172716  
 -.172716      0.931470

\*\*\*\*\*

## R E P A R T: T T M A T R I X

-----  
 0.600459      0.317924E-01  
 -.317924E-01   -.276468

\*\*\*\*\*

## I M P A R T: T T M A T R I X

-----  
 -.729781      -.150166  
 0.150166      0.134581E-01

\*\*\*\*\*

## Z      R E P A R T: F U N C T I O N S

-----  
 -25.0000    0.2716D+00    0.1259D+00    0.1631D+00    -.5370D+00    0.1506D-05    -.4406D-04    0.6915D-07    -.2284D-05  
 -13.6000    0.5769D+00    0.6829D-01    0.8088D-01    -.1298D+01    -.3777D-04    0.5932D-04    -.3999D-05    0.9632D-05  
 -3.6000    -.4115D+00    0.6691D-01    -.8351D-01    0.1351D+01    -.4630D-02    -.2048D-01    0.2890D-04    -.9308D-04  
 0.0000    -.8416D+00    -.7861D-01    -.9335D-02    0.4446D+00    0.5115D-01    0.1732D-01    0.2247D-02    -.6850D-02  
 3.6000    0.1490D-01    0.5461D-01    0.3718D-01    -.2780D+00    -.9230D-02    -.2404D-02    0.1299D-02    -.2425D-03  
 13.6000    0.6802D+00    0.7978D-01    -.4223D-01    0.2431D+00    -.2701D-04    -.3867D-04    0.2948D-05    0.3128D-05  
 25.0000    0.6664D+00    0.1165D+00    -.1531D+00    -.1120D+00    0.7601D-06    -.8680D-05    -.2445D-07    0.4751D-06

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## Z      I M P A R T: F U N C T I O N S

-----  
 -25.0000    -.8982D+00    -.3837D-01    -.6149D-01    -.6546D+00    0.5103D-05    -.4498D-04    0.2851D-06    -.2320D-05  
 -13.6000    -.5507D+00    -.1129D+00    0.1559D+00    -.1335D+01    0.6059D-05    0.8405D-04    0.3894D-06    0.1222D-04  
 -3.6000    -.5284D+00    0.8131D-01    -.1780D+00    0.1380D+01    0.1289D-01    -.1938D-01    0.2320D-02    0.4662D-04  
 0.0000    0.5262D+00    0.1487D+00    0.1846D-01    0.6235D+00    -.3508D-01    0.4223D-03    0.5965D-02    -.9388D-02  
 3.6000    0.7372D+00    0.1083D+00    0.1518D+00    0.1221D+00    0.1107D-02    0.5541D-02    0.1592D-02    -.3799D-03  
 13.6000    -.2428D+00    -.8603D-01    -.1506D+00    -.1425D+00    0.2784D-04    -.5248D-04    -.2187D-05    0.5597D-05  
 25.0000    0.2735D+00    -.1055D-01    0.2391D-01    -.2560D+00    0.6563D-05    0.4645D-05    -.3403D-06    -.2162D-06

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## C H E C K   P R O P E R T I E S

```

-----
*****
          C H E C K |RR_<-|^2 + |TT_<-|^2
-----
    1.00000    0.242339E-09
    0.242339E-09    1.00000

*****

          C H E C K |RR_->|^2 + |TT_->|^2
-----
    1.00000    -.407011E-09
   -.407011E-09    1.00000

*****

          R E P A R T: TT_->^1 * RR_<- + RR_->^1 * TT_<-
-----
    0.185469E-09  0.420999E-09
   -.476236E-09  0.157399E-09

          I M P A R T: TT_->^1 * RR_<- + RR_->^1 * TT_<-
-----
    0.219235E-11 -.125379E-09
   -.197244E-09  0.129723E-10

*****

          R E P A R T: RR_<-^T - RR_<-
-----
    0.00000    -.185546E-09
    0.185546E-09    0.00000

          I M P A R T: RR_<-^T - RR_<-
-----
    0.00000    0.356981E-09
   -.356981E-09    0.00000

*****

          R E P A R T: RR_->^T - RR_->
-----
    0.00000    0.103188E-09
   -.103188E-09    0.00000

          I M P A R T: RR_->^T - RR_->
-----
    0.00000    -.533526E-09
    0.533526E-09    0.00000

*****

          R E P A R T: TT_->^T - TT_<-
-----
    0.231348E-10  0.847061E-10
   -.952086E-13  0.142655E-10

          I M P A R T: TT_->^T - TT_<-
-----
    0.186473E-11  0.452252E-09
    0.511466E-09 -.116038E-10
*****

```

## References

1. A Program Package for Solution of Two-Dimensional Discrete and Continuum Spectra Boundary-Value Problems in Kantorovich (Adiabatic) Approach / O. Chuluunbaatar, A. A. Gusev, S. I. Vinitzky, A. G. Abrashkevich // JINR Lib. — 2013. — <http://wwwinfo.jinr.ru/programs/jinrlib/kantbp/indexe.html>.
2. Symbolic-Numerical Algorithms to Solve the Quantum Tunneling Problem for a Coupled Pair of Ions / A. A. Gusev, S. I. Vinitzky, O. Chuluunbaatar et al. // Lecture Notes in Computer Science. — 2011. — Vol. 6885. — Pp. 175–191.
3. KANTBP: A Program for Computing Energy Levels, Reaction Matrix and Radial Wave Functions in the Coupled-Channel Hyperspherical Adiabatic Approach / O. Chuluunbaatar, A. A. Gusev, A. G. Abrashkevich et al. // Comput. Phys. Commun. — 2007. — Vol. 177. — Pp. 649–675.
4. KANTBP 2.0: New Version of a Program for Computing Energy Levels, Reaction Matrix and Radial Wave Functions in the Coupled-Channel Hyperspherical Adiabatic Approach / O. Chuluunbaatar, A. A. Gusev, S. I. Vinitzky, A. G. Abrashkevich // Comput. Phys. Commun. — 2008. — Vol. 179. — Pp. 685–693.
5. Symbolic-Numerical Algorithm for Generating Cluster Eigenfunctions: Tunneling of Clusters Through Repulsive Barriers / A. A. Gusev, S. I. Vinitzky, O. Chuluunbaatar et al. // Lecture Notes in Computer Science. — 2013. — Vol. 8136. — Pp. 427–442.
6. ODPEVP: A program for Computing Eigenvalues and Eigenfunctions and Their First Derivatives with Respect to the Parameter of the Parametric Self-Adjoined Sturm–Liouville Problem / O. Chuluunbaatar, A. A. Gusev, S. I. Vinitzky, A. G. Abrashkevich // Comput. Phys. Commun. — 2009. — Vol. 180, No 8. — Pp. 1358–1375.

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### **KANTBP 3.0: новая версия программы для вычисления энергетических уровней, матриц амплитуд отражения и прохождения и соответствующих волновых функций в адиабатическом подходе со связанными каналами**

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Представлено краткое описание программ на языке Фортран 77 для вычисления энергетических уровней, матриц амплитуд отражения и прохождения и соответствующих волновых функций в адиабатическом подходе со связанными каналами. В этом подходе многомерное уравнение Шрёдингера сводится к системе связанных обыкновенных дифференциальных уравнений второго порядка на конечном интервале с однородными граничными условиями третьего рода на левой и правой граничных точках для задачи непрерывного спектра или набора граничных условий первого, второго и третьего рода для задачи дискретного спектра. Полученная система уравнений, содержащая матричные потенциалы, а также связанная слагаемыми, содержащими первые производные, решается в приближении высокого порядка точности методом конечных элементов.

**Ключевые слова:** краевая задача, многоканальная задача рассеяния, метод конечных элементов, метод Канторовича.